Chapter 12. The Chemical Bond in Diatomic Molecules

I. Born-Oppenheimer approximation

Because of the small mass of an electron, it can respond instantaneously to change of nuclear coordinates.

$$\Psi(\mathbf{r},\mathbf{R}) = \psi(\mathbf{r}:\mathbf{R})\varphi(\mathbf{R})$$

Substituting to $\hat{H}\Psi = E\Psi$ and ignoring the action of \hat{T}_n on $\psi(\mathbf{r}: \mathbf{R})$, we have

$$[\hat{T}_{e} + V_{en}]\psi(\mathbf{r}:\mathbf{R}) = E_{e}(\mathbf{R})\psi(\mathbf{r}:\mathbf{R}) \qquad (\text{electronic})$$
$$[\hat{T}_{n} + V_{nn} + E_{e}(\mathbf{R})]\varphi(\mathbf{R}) = E\varphi(\mathbf{R}) \qquad (\text{nuclear})$$

In B-O approximation, electronic Schrödinger equation is solved at a fixed nuclear framework. The nuclei Schrödinger equation is solved on <u>potential energy surface</u> $(V = V_{nn} + E_e)$.

Molecular Orbital Theory

Hydrogen molecular ion

H₂⁺ electronic Hamiltonian

$$\hat{H}_{e} = \hat{T}_{e} + V_{en} = -\frac{\hbar^{2}}{2m_{e}}\nabla_{e}^{2} + \frac{e^{2}}{4\pi\varepsilon_{0}}\left(-\frac{1}{r_{a}} - \frac{1}{r_{b}}\right)$$

LCAO (linear combination of atomic orbitals):

$$\Psi^{MO} = \sum_{n} c_n \psi^{AO}$$

Let's assume that only the 1s orbitals are involved:

$$\Psi = c_a 1 s_a + c_b 1 s_b$$

At fixed R, Schrödinger Eq. for $\hat{H}_e = \hat{T}_e + V_{en}$

$$\hat{H}_e \Psi = E \Psi$$

Substituting LCAO-MO, we have

$$c_a \hat{H}_e 1 s_a + c_b \hat{H}_e 1 s_b = c_a E 1 s_a + c_b E 1 s_b$$

Multiply $1s_a$ on the left and integrate:

$$H_{aa}c_a + H_{ab}c_b = Ec_a + S_{ab}Ec_b$$

We define the <u>Coulomb integrals</u>:

$$H_{aa} = \int 1s_a \hat{H}_e 1s_a d\tau , \qquad H_{bb} = \int 1s_b \hat{H}_e 1s_b d\tau$$

the resonance integral:

$$H_{ab} = \int 1s_a \hat{H}_e 1s_b d\tau = H_{ba}$$

and the overlap integral:

$$S_{ab} = \int 1s_a 1s_b d\tau = S_{ba}$$

Similarly with $1s_b$,

$$H_{ab}c_a + H_{bb}c_b = ES_{ab}c_a + Ec_b$$

Matrix form of the secular equation:

$$\begin{pmatrix} H_{aa} - E & H_{ab} - ES_{ab} \\ H_{ab} - ES_{ab} & H_{bb} - E \end{pmatrix} \begin{pmatrix} c_a \\ c_b \end{pmatrix} = 0$$

Solution of coupled linear equations requires the <u>secular</u> <u>determinant</u> to be zero:

$$\begin{vmatrix} H_{aa} - E & H_{ab} - ES_{ab} \\ H_{ab} - ES_{ab} & H_{bb} - E \end{vmatrix} = 0$$

For homonuclear diatoms, $H_{aa} = H_{bb} = \alpha$, $H_{ab} = \beta$, $S_{ab} = S$

$$(\alpha - E)^2 - (\beta - ES)^2 = 0$$

Solution:

$$E_{\pm} = \frac{\alpha \pm \beta}{1 \pm S}$$

Bonding and antibonding MOs

$$\Psi_{-}$$
 1s_a 1s_b Ψ_{+}

Substituting energy back to secular equation, we have

$$c_a^+ = c_b^+ = \frac{1}{\sqrt{2(1+S)}}, \qquad c_a^- = -c_b^- = \frac{1}{\sqrt{2(1-S)}}$$

Wavefunctions

$$\Psi_{+} = \frac{1}{\sqrt{2(1+S)}} (1s_a + 1s_b), \qquad \Psi_{-} = \frac{1}{\sqrt{2(1-S)}} (1s_a - 1s_b)$$

$$P_{+} = N_{+}^{2} [1s_{a}^{2} + 1s_{b}^{2} + 2(1s_{a}1s_{b})]$$

<u>Bonding MO</u> has extra e-density between nuclei (constructive interference). It lowers the energy of MO relative to AOs.

$$P_{-} = N_{-}^{2} [1s_{a}^{2} + 1s_{b}^{2} - 2(1s_{a}1s_{b})]$$

<u>Antibonding MO</u> has a node (destructive interference). It raises the energy of MO relative to AOs.

Chemical bond is possible because of the interference, thus a quantum phenomenon.

Dependence on R (potential energy curves)

$$R_e = 1.3 \text{\AA}$$

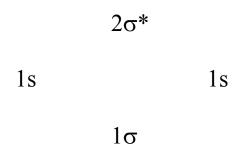
 $D_e = 1.77 \text{ eV}$

Homonuclear diatoms

H₂:

MOs are essentially the same as in H_2^+ .

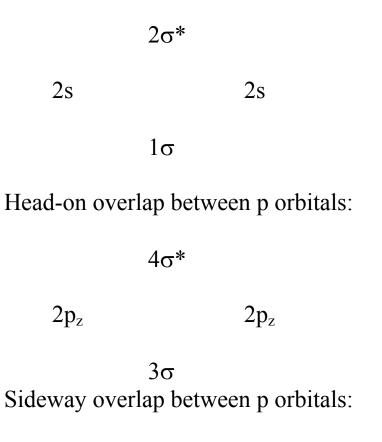
MOs formed by s-type AOs are called σ orbitals, because they are cylindrically symmetric about internuclear axis with zero angular momentum.



Both electrons are in the bonding orbital.

2nd period homonuclear diatoms:

1s electrons are in the core and do not interact strongly. Valence electrons:



 $2\pi^*$

 1π

 π MOs have 1 unit of angular momentum around the axis. Both π MOs are doubly degenerate.

MO energy level diagram

Electronic configuration determined by <u>Aufbau principle</u>: H_2^+ , $1\sigma^1$

H₂, 1σ², spin paired (Pauli principle), Bond order = (# bonding e⁻ - # antibonding e⁻)/2 = (2-0)/2 = 1, single bond.

He₂, $1\sigma^2 2\sigma^{*2}$, unstable, energy gained in σ offset by energy loss in σ^* , bond order = 0

 Li_2 , $1\sigma^2$, single bond.

HOMO: highest occupied MO (1σ) LUMO: lowest unoccupied MO $(2\sigma*)$.

These orbitals are called <u>frontier orbitals</u> and largely responsible for chemical and spectroscopic properties of the molecule.

 $N_{2}, 1\sigma^{2}2\sigma^{*2}1\pi^{4}3\sigma^{2},$

Extremely stable because of 6 bonding e- (triple bond).

Because of interaction with $2\sigma^*$, 3σ is higher than 1π . This is confirmed by photoelectron spectrum.

Relative bond strength:

N₂, $1\sigma^2 2\sigma^{*2} 1\pi^4 3\sigma^2$, bond order = (8-2)/2=3 N₂⁺, $1\sigma^2 2\sigma^{*2} 1\pi^4 3\sigma^1$, bond order = (7-2)/2=2.5 N₂⁻, $1\sigma^2 2\sigma^{*2} 1\pi^4 3\sigma^2 2\pi^{*1}$, bond order = (8-3)/2=2.5 O₂, $1\sigma^2 2\sigma^{*2} 3\sigma^2 1\pi^4 2\pi^{*2}$, bond order = 2, two top electrons unpaired (Hund's rule), paramagnetic.

F₂,
$$1\sigma^2 2\sigma^{*2} 3\sigma^2 1\pi^4 2\pi^{*4}$$
, bond order = 1.
Ne₂, $1\sigma^2 2\sigma^{*2} 3\sigma^2 1\pi^4 2\pi^{*2} 4\sigma^{*2}$, bond order = 0, unstable.

Heteronuclear diatoms

Secular determinant:

$$\begin{vmatrix} H_{aa} - E & H_{ab} - ES_{ab} \\ H_{ab} - ES_{ab} & H_{bb} - E \end{vmatrix} = 0$$

Assuming:

$$H_{aa} = \alpha_a, \qquad H_{bb} = \alpha_b, \qquad H_{ab} = \beta, \qquad S_{ab} = 0$$

Solution can be expressed in terms of ζ :

$$E_{+} = \alpha_{a} + \beta \tan \zeta$$
$$E_{-} = \alpha_{b} - \beta \tan \zeta$$

where

$$\tan 2\zeta = \frac{2|\beta|}{\alpha_b - \alpha_a}$$

 ΔE of two AOs $(\alpha_b - \alpha_a)$ determines the strength of MO. Example: H-F bond

Energies (α) of AOs: H1s: -13.6 eV, F2p: -18.6 eV

Assume $\beta = -1.0 \text{ eV}$, $\zeta = 10.9^{\circ}$,

$$E_{+} = -18.6 - 1.0 \tan(10.9^{\circ}) = -18.8 \text{ eV}$$

$$E_{-} = -13.6 + 1.0 \tan(10.9^{\circ}) = -13.4 \text{ eV}$$

CO:

NO:

 $1\sigma^2 2\sigma^{*2} 1\pi^4 3\sigma^2$ Triple bond, very stable. $1\sigma^2 2\sigma^{*2} 1\pi^4 3\sigma^2 \pi^{*1}$ Free radical