

The quaternionic contact Yamabe problem on a  
3-Sasakian manifold.

Dimitar Vassilev, University of New Mexico

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## The equations

- ▶ Yamabe:  $(M^n, g)$ :  $\bar{g} = u^{4/(n-2)} g$ ,

$$4 \frac{n-1}{n-2} \Delta u - S u = -\bar{S} u^{2^*-1}.$$

- ▶ CR Yamabe:  $(M^{2n+1}, \eta, J)$ ,  $\bar{\eta} = u^{4/(Q-2)} \eta$ ,

$$4 \frac{n+1}{n} \Delta u - S u = -\bar{S} u^{2^*-1}.$$

- ▶ QC Yamabe:  $(M^{4n+3}, \eta)$ ,  $\bar{\eta} = u^{4/(Q-2)} \eta$ ,

$$4 \frac{n+2}{n+1} \Delta u - S u = -\bar{S} u^{2^*-1}.$$

- ▶  $\Delta f = -\lambda_1 f$ .

## Riemannian Obata theorems

**Theorem 1.** a) (Uniqueness in Einstein class) Let  $(M, \bar{g})$  be a connected compact Riemannian manifold. If  $\bar{g}$  is Einstein and  $g = \phi^2 \bar{g}$  with  $\bar{S} = S = n(n-1)$ , then  $\phi = 1$  unless  $(M, \bar{g})$  is the round unit sphere  $(S^n, g_{st})$ .

b) (Yamabe problem on the round sphere) If  $g$  is conformal to  $g_{st}$  on  $S^n$ ,  $g = \phi^2 g_{st}$ , with  $S = n(n-1)$ , then  $g = \Phi^* g_{st}$  for  $\Phi \in \text{Diff}(S^n)$ .

**Theorem 2.** Let  $(M, g)$  be an  $n$ -dimensional compact Riemannian manifold with

$$\text{Ric}(X, X) \geq (n-1)g(X, X).$$

If  $\lambda \neq 0$  is an eigenvalue,  $\Delta f = -\lambda f$ , then  $\lambda \geq n$  (Lichnerowicz) and  $\lambda = n$  iff  $(M, g)$  is isometric with  $S^n$  (Obata), in which case  $f$  is a spherical harmonic of order one.

## The PDE on $\mathbb{R}^n$ - extremals of the $L^2$ Sobolev embedding inequality

Stereographic proj.,  $\mathcal{C} : S^n \setminus N \rightarrow \mathbb{R}^n$ ,  $(\mathcal{C}^{-1})^* g_{st} = 4u^{4/(n-2)} dx^2$ . The Yamabe problem on the round sphere is equivalent to:

**Theorem 3.** (Aubin, Talenti) If  $u \geq 0$  satisfies the Yamabe equation on  $\mathbb{R}^n$

$$\Delta u = -n(n-2)u^{2^*-1}, \quad u \in \mathcal{D}^{1,2}(\mathbb{R}^n)$$

then up to a translation and rescaling  $u = (1 + |x|^2)^{-(n-2)/2}$ .

$$\text{Rescaling: } u_\lambda(x) \equiv \lambda^{n/2^*} \delta_{\lambda u} \stackrel{\text{def}}{=} \lambda^{n/2^*} u(\lambda x), \quad \lambda > 0.$$

Key: reduce to radial functions via symmetrization arguments. These are not (fully) available in sub-Riemannian settings (ex. groups of Iwasawa type) except for solutions with "partial" symmetry w/ Garofalo or the lowest energy solutions (extremals for Folland-Stein  $L^2$  Sobolev type inequality): Branson & Fontana & Morpurgo and Frank & Lieb in the CR case, w/ Ivanov - Minchev in the quaternion case, Christ & Liu & Zhang in the octonian case.

**Uniqueness.** Recall,  $g = \phi^2 \bar{g}$  and  $S = \bar{S} = n(n-1)$

Suppose  $\bar{g}$  is Einstein,  $0 = \overline{Ric}_o = Ric_o + \frac{n-2}{\phi}(\nabla^2\phi)_0$ . The contracted Bianchi identity and  $S=\text{const}$  give  $\nabla^* Ric_0 = \frac{n-2}{2n}\nabla S = 0$ , hence

$$\nabla^* (Ric_o \nabla\phi) = (\nabla^* Ric_o)(\nabla\phi) + g(Ric_o, \nabla^2\phi) = \frac{n-2}{2n}g(\nabla S, \nabla\phi) - \frac{\phi}{n-2}|Ric_o|^2.$$

This *divergence formula* shows that  $g$  is *also an Einstein metric* and  $X = \nabla\phi$  is a *gradient conformal vector field*,

$$Ric_o = (\nabla^2\phi)_0 = 0.$$

If  $X$  is a conformal vector field then we have the *infinitesimal Yamabe equation*

$$\Delta(\text{div } X) = -\frac{1}{n-1}(\text{div } X)S - \frac{n}{2(n-1)}X(S).$$

Now, for  $S = n(n-1)$  it follows  $f = \Delta\phi$  satisfies  $\Delta f = -nf$ . Thus either  $f = \text{const}$  or  $f$  is an eigenfunction with the lowest possible eigenvalue hence  $g$  is isometric to  $g_{st}$  by *Obata's eigenvalue theorem*.

## The case of the sphere

Taking into account *the divergence formula*, using the stereographic projection we can reduce to a conformal map of the Euclidean space, which sends the *Euclidean metric to a conformal to it Einstein metric*. By a purely local argument *the resulting system can be integrated*, in effect proving also *Liouville's theorem*, which gives the form of  $u$  as in Aubin and Talenti's theorem in  $\mathbb{R}^n$  and then  $\phi$  on  $S^n$  after transferring the equations back to the unit sphere.

Remark: Such argument was used in the quaternionic contact setting to classify all qc-Einstein structures on the unit  $4n + 3$  dimensional sphere (quaternionic Heisenberg group) conformal to the standard qc-structure.

## OBATA TYPE RESULTS ON CR AND QC MANIFOLDS

## Sub-Riemannian conformal infinities

On the open unit ball  $B$  in  $\mathbb{C}^{n+1}$  consider the Bergman metric

$$h = \frac{4}{\rho} g_{euc} + \frac{1}{\rho^2} ((d\rho)^2 + (Id\rho)^2), \quad \rho = 1 - |x|^2.$$

As  $\rho \rightarrow 0$ ,  $\rho \cdot h$  is finite only on  $H = \text{Ker}(Id\rho)$ , which is the kernel of the contact form  $\theta = Id\rho$ . The conformal infinity of  $\rho \cdot h$  is the conformal class of a pseudohermitian CR structure on  $S^{2n+1}$ .

In the quaternion case, consider the open unit ball  $B$  in  $\mathbb{H}^{n+1}$  and the hyperbolic metric

$$h = \frac{4}{\rho} g_{euc} + \frac{1}{\rho^2} ((d\rho)^2 + (I_1 d\rho)^2 + (I_2 d\rho)^2 + (I_3 d\rho)^2).$$

The conformal infinity is the conformal class of a (QC) quaternionic contact structure on  $S^{4n+3}$ . Here,  $\rho \cdot h$  defines a conformal class of degenerate metrics with kernel

$$H = \bigcap_{j=1}^3 \text{Ker}(I_j d\rho),$$

which carries a quaternionic structure.



## CR SETTING

## The solution in the Sasaki-Einstein case

### ***Theorem 4.***

- a) (*Jerison & Lee '88*) If  $\theta$  is the contact form of a pseudo-Hermitian structure proportional to the standard contact form  $\bar{\theta}$  on the unit sphere in  $\mathbb{C}^{n+1}$  and the pseudohermitian scalar curvature  $S_\theta = \text{const}$ , then up to a multiplicative constant  $\theta = \Phi^* \bar{\theta}$  with  $\Phi$  a CR automorphism of the sphere.
- b) (*X. Wang '13, Ivanov & Vassilev '14*) The pseudoconformal class of a Sasaki-Einstein pseudo-Hermitian structure different from the standard Sasaki-Einstein structure on the round sphere contains a unique (up to homothety) pseudo-Hermitian form of constant CR scalar curvature.

## CR manifolds

$(M, \theta, J)$  is **strictly pseudoconvex pseudo-Hermitian manifold** if

- i)  $\theta$  is a contact form,  $H = \ker \theta$  has a compatible Hermitian structure:  $J : H \rightarrow H, J^2 = -id_H, 2g(X, Y) \stackrel{def}{=} d\theta(X, JY), X, Y \in H, g(X, Y) = g(JX, JY)$ ;
- ii)  $g$  is positive definite on  $H$ ;
- iii) integrability:  $[JX, Y] + [X, JY] \in H$  and  $[JX, JY] - [X, Y] - J[JX, Y] - J[X, JY] = 0$ .

**Reeb field**  $\xi$ :  $\theta(\xi) = 1$  and  $\xi \lrcorner d\theta = 0$ .

**Tanaka-Webster connection.** Unique linear connection  $\nabla$  such that

(i)  $\xi, J, \theta$  and  $g$  are parallel; (ii) the torsion satisfies:

- ▶  $T(X, Y) = 2\omega(X, Y)\xi$ , where  $\omega(X, Y) \stackrel{def}{=} g(JX, Y), X, Y \in H$ ;
- ▶ the Webster torsion  $A, A \stackrel{def}{=} T(\xi, \cdot) : H \rightarrow H$ , is symmetric and anti-commutes with  $J, AJ = -JA$ .

Note:  $A = 0 \Leftrightarrow$  Sasakian structure  $\Leftrightarrow \mathcal{L}_\xi g = 0$ .

## Curvature of the Tanaka-Webster connection

Define the Riemannian metric " $h = g + \eta^2$ ". Let  $\{\epsilon_a\}_{a=1}^{2n}$ -ONB of the horizontal space  $H$ .

- ▶ Tanaka-Webster curvature:  $R(A, B)C \stackrel{\text{def}}{=} [\nabla_A, \nabla_B]C - \nabla_{[A, B]}C$   
and  $R(A, B, C, D) \stackrel{\text{def}}{=} h(R(A, B)C, D)$ .
- ▶ Ricci tensor:  $Ric(A, B) = R(\epsilon_a, A, B, \epsilon_a) \stackrel{\text{def}}{=} \sum_{a=1}^{2n} R(\epsilon_a, A, B, \epsilon_a)$ ;  
scalar curvature  $S = Ric(\epsilon_a, \epsilon_a)$ ;
- ▶ Ricci form:  $\rho(A, B) = \frac{1}{2} R(A, B, \epsilon_a, J\epsilon_a)$ .

**Proposition 5.** *We have the following type decomposition of the Ricci tensor with  $B = \rho_0$ ,  $\rho(JX, Y) = B(JX, Y) + \frac{1}{2n}g(X, Y)$ ,*

$$Ric(X, Y) = 2(n-1)A(JX, Y) + B(JX, Y) + \frac{S}{2n}g(X, Y).$$

A torsion-free pseudo-Einstein CR manifold is Sasaki-Einstein if  $S = 4n(n+1)$ .

## CR divergence formula

**Theorem 6 (Jerison, D. & Lee, J. '88).** *Let  $(M, \bar{\theta})$  be a compact Sasaki-Einstein manifold. If  $\theta = 2h\bar{\theta}$  is also of constant positive pseudo-Hermitian scalar curvature  $S = 4n(n + 1)$ , then  $(M, \theta)$  is again a Sasaki-Einstein space.*

”Proof”: Divergence formula: for a certain horizontal vector field  $X_h$  we have

$$\nabla^* X_h = \frac{1}{2} \left( \frac{1}{2} + h \right) (|D|^2 + |E|^2) + \frac{h}{4} |D^h + E^h|^2 + \frac{h}{2} Q(d, e, u),$$

where  $D(X, Y) = -4A(X, Y)$  and  $E = \frac{2}{n+2}B(X, Y)$  are up to a constant multiple the Webster torsion and the traceless  $J$ -invariant component of the Ricci tensor of the Tanaka-Webster connection.

With  $f = \frac{1}{2} + h + \frac{|\nabla h|^2}{4h}$ ,  $d = h^{-1}DJ\nabla h$ ,  $e = h^{-1}EJ\nabla h$ , and  $u = \frac{1}{n+2}\nabla^*(JD)$  we have

$$X_h = f[d + e] - dh(\xi)(Jd - Je + 6Ju).$$

**Infinitesimal CR transformations:**  $\mathcal{L}_Q \theta = f \theta$  and  $\mathcal{L}_Q J = 0$ .

If  $\bar{\theta} = \Phi_t^* \theta = u_t^{2/n} \theta$ , then  $\frac{4(n+1)}{n} \Delta u_t - S u_t = -(S \circ \Phi_t) u_t^{2^*-1}$ .

Differentiating at  $t = 0$ , the function  $\phi = \frac{d}{dt} u_t|_{t=0}$  satisfies

$$\frac{4(n+1)}{n} \Delta \phi - S \phi = -dS(Q) - S(2^* - 1) \phi.$$

**Proposition 7.** *An infinitesimal CR automorphism  $Q$  satisfies*

$$\Delta(\nabla^* Q_H) = -\frac{n}{2(n+1)} dS(Q) - \frac{S}{2(n+1)} \nabla^* Q_H.$$

Proof: Use  $\mathcal{L}_Q g(X, Y) = \frac{1}{n} (\nabla^* Q_H) g(X, Y)$ . Hence  $\bar{g} = u_t^{2/n} g$  gives  $\frac{2}{n} \phi = \frac{1}{n} \nabla^* Q_H$ . Also  $2^* = \frac{2(n+1)}{n}$ .

**Characterization:**

- ▶  $f = d\sigma(\xi)$  and  $Q = -\frac{1}{2} J \nabla \sigma - \sigma \xi$ , where  $Q_H$  ("contact Hamiltonian field") is determined by  $\theta(Q_H) = 0$  and  $i_{Q_H} d\theta \equiv 0 \pmod{\theta}$ ;
- ▶  $[\nabla^2 \sigma]_{[-1]}(X, Y) \equiv \frac{1}{2} [\nabla^2 \sigma(X, Y) - \nabla^2 \sigma(JX, JY)] = -2\sigma A(JX, Y)$ .

## Consequences of $\Delta(\nabla^* Q_H) = -\frac{n}{2(n+1)}dS(Q) - \frac{S}{2(n+1)}\nabla^* Q_H$

When  $A = 0$ , Ricci's identity gives  $\nabla^3 h(X, Y, \xi) = \nabla^3 h(\xi, X, Y)$  while  $\theta = 2h\bar{\theta}$  gives  $[\nabla^2 h]_{[-1]}(X, Y) = -2hA(X, JY) = 0$ . Hence, the vector field

$$Q = -\frac{1}{2}J\nabla(\xi h) - (\xi h)\xi$$

is an infinitesimal CR vector field unless it vanishes. Since  $S = 4n(n+1)$  it follows  $\phi = \nabla^* Q_H$  either vanishes identically, i.e.,  $h = \text{const}$  or  $\phi$  is an eigenfunction of the sublaplacian realizing the smallest possible eigenvalue on a (pseudo-Einstein) Sasakian manifold and  $h \neq \text{const}$ .

The CR Lichnerowicz-Obata theorem shows that  $(M, \theta)$  is homothetic to the CR unit sphere.

Remark: For  $f = \frac{1}{2} + h + \frac{|\nabla h|^2}{4h}$  it follows  $Q = -\frac{1}{2}\nabla f - dh(\xi)\xi$  and  $\phi = \Delta f$ .

## CR Lichnerowicz theorem

*Theorem 8 (Greenleaf, A. '85 for  $n \geq 3$ ; Li, S.-Y., & Luk, H.-S. '04 for  $n=2$ ). Let  $M$  be a compact spcph manifold of dimension  $2n + 1$ , s.t., for some  $k_0 = \text{const} > 0$  we have the Lichnerowicz-type bound*

$$\text{Ric}(X, X) + 4A(X, JX) \geq k_0 g(X, X), \quad X \in H.$$

*If  $n > 1$ , then any eigenvalue  $\lambda$  of the sub-Laplacian satisfies*

$$\lambda \geq \frac{n}{n+1} k_0.$$

The standard Sasakian unit sphere has first eigenvalue equal to  $2n$  with eigenspace spanned by the restrictions of all linear functions to the sphere.



**Theorem 9 (Chiu, H.-L. '06).** *If  $n = 1$  the estimate  $\lambda \geq \frac{n}{n+1}k_0$  holds assuming in addition that the CR-Paneitz operator is non-negative  $\int_M f \cdot Cf \text{Vol}_\theta \geq 0$ , where  $Cf$  is the CR-Paneitz operator,*

$$Cf = \nabla^4 f(e_a, e_a, e_b, e_b) + \nabla^4 f(e_a, Je_a, e_b, Je_b) \\ - 4n \nabla^* A(J\nabla f) - 4n g(\nabla^2 f, JA).$$

Note: Li, S.-Y., & Luk, H.-S. '04 for  $n = 1$  with condition.

Given a function  $f$  we define the one form,

$$P_f(X) = \nabla^3 f(X, e_b, e_b) + \nabla^3 f(JX, e_b, Je_b) + 4nA(X, J\nabla f)$$

so we have  $Cf = -\nabla^* P$ .

## CR Obata type theorem

**Theorem 10** ( $n \geq 2$ , Li, S.-Y., Wang, X. '13;  $n=1$  w/ Ivanov '14).  
Suppose  $(M, J, \theta)$ ,  $\dim M = 2n + 1$ , is a compact spcph manifold which satisfies the Lichnerowicz-type bound. If  $n \geq 2$ , then  $\lambda = \frac{n}{n+1}k_0$  is an eigenvalue iff up-to a scaling  $(M, J, \theta)$  is the standard pseudo-Hermitian CR structure on the unit sphere in  $\mathbb{C}^{n+1}$ . If  $n = 1$  the same conclusion holds assuming in addition  $C \geq 0$ .

### Earlier results

- ▶ Sasakian case (enough for the CR Yamabe problem on the sphere!), Chang, S.-C., & Chiu, H.-L., for  $n \geq 2$  in J. Geom. Anal. '09; for  $n = 1$  in Math. Ann. '09.
- ▶ Non-Sasakian case, Chang, S.-C., & Wu, C.-T., '12, assuming:  
(i) for  $n \geq 2$ ,  $A_{\alpha\beta, \bar{\beta}} = 0$  and  $A_{\alpha\beta, \gamma\bar{\gamma}} = 0$ ; (ii) for  $n = 1$ ,  $A_{11, \bar{1}} = 0$  and  $P_1 f = 0$ .
- ▶ w/ Ivanov '12 - assuming  $\nabla^* A = 0$  and  $C \geq 0$  when  $n = 1$ .

## QUATERNIONIC CONTACT CASE

## Solution of the Yamabe problem in the 3-Sasakian case

**Theorem 11 (w/ Ivanov & Minchev arXiv:1504.03142).** a) Let  $(M, \bar{\eta})$  be a compact locally 3-Sasakian qc manifold of qc-scalar curvature  $16n(n+2)$ . If  $\eta = 2h\bar{\eta}$  is qc-conformal to  $\bar{\eta}$  structure which is also of constant qc-scalar curvature, then up to a homothety  $(M, \eta)$  is locally 3-Sasakian manifold. Furthermore, the function  $h$  is constant unless  $(M, \bar{\eta})$  is the unit 3-Sasakian sphere.

b) Let  $\eta = 2h\tilde{\eta}$  with  $\tilde{\eta}$  the standard qc-structure on a 3-Sasakian sphere of dimension  $4n+3$ . If  $\eta$  has constant qc-scalar curvature  $16n(n+2)$ , then  $\eta$  is obtained from  $\tilde{\eta}$  by a conformal quaternionic contact automorphism.

Remark: The 7-D case of b) was completed earlier '10.

The qc-Yamabe equation on the quaternionic Heisenberg group of homogeneous dimension  $Q = 4n + 6$

**Corrolary 12.** *If  $0 \leq \Phi \in \mathcal{D}^{1,2}(\mathbf{G}(\mathbb{H}))$ ,  $S_\Theta = \text{const}$ ,*

$$\frac{4(Q+2)}{Q-2} \Delta_{\tilde{\Theta}} \Phi = -S_\Theta \Phi^{2^*-1},$$

*then for some fixed  $(q_o, \omega_o) \in \mathbf{G}(\mathbb{H})$ , constants  $c_0 > 0$  and  $\sigma > 0$  such that  $S_\Theta = 128n(n+2)c_0\sigma$  we have  $\Phi = (2h)^{-(Q-2)/4}$  with*

$$h(q, \omega) = c_0 \left[ (\sigma + |q + q_o|^2)^2 + |\omega + \omega_o + 2 \operatorname{Im} q_o \bar{q}|^2 \right].$$

The sub-Laplacian is  $\Delta_{\tilde{\Theta}} u = \sum_{\alpha=1}^n (T_\alpha^2 u + X_\alpha^2 u + Y_\alpha^2 u + Z_\alpha^2 u)$ .

# Quaternionic Heisenberg Group

$$\mathbf{G}(\mathbb{H}) = \mathbb{H}^n \times \text{Im}\mathbb{H}, \quad (q, \omega) \in \mathbf{G}(\mathbb{H}),$$

$$(q_o, \omega_o) \circ (q, \omega) = (q_o + q, \omega + \omega_o + 2 \text{Im } q_o \bar{q}),$$

i)  $\tilde{\Theta} = (\tilde{\Theta}_1, \tilde{\Theta}_2, \tilde{\Theta}_3) = \frac{1}{2} (d\omega - q \cdot d\bar{q} + dq \cdot \bar{q})$  or

$$\tilde{\Theta}_1 = \frac{1}{2} dx - x^\alpha dt^\alpha + t^\alpha dx^\alpha - z^\alpha dy^\alpha + y^\alpha dz^\alpha$$

$$\tilde{\Theta}_2 = \frac{1}{2} dy - y^\alpha dt^\alpha + z^\alpha dx^\alpha + t^\alpha dy^\alpha - x^\alpha dz^\alpha$$

$$\tilde{\Theta}_3 = \frac{1}{2} dz - z^\alpha dt^\alpha - y^\alpha dx^\alpha + x^\alpha dy^\alpha + t^\alpha dz^\alpha.$$

ii) Left-invariant horizontal vector fields

$$T_\alpha = \frac{\partial}{\partial t_\alpha} + 2x^\alpha \frac{\partial}{\partial x} + 2y^\alpha \frac{\partial}{\partial y} + 2z^\alpha \frac{\partial}{\partial z}, \quad X_\alpha = \frac{\partial}{\partial x_\alpha} - 2t^\alpha \frac{\partial}{\partial x} - 2z^\alpha \frac{\partial}{\partial y} + 2y^\alpha \frac{\partial}{\partial z},$$

$$Y_\alpha = \frac{\partial}{\partial y_\alpha} + 2z^\alpha \frac{\partial}{\partial x} - 2t^\alpha \frac{\partial}{\partial y} - 2x^\alpha \frac{\partial}{\partial z}, \quad Z_\alpha = \frac{\partial}{\partial z_\alpha} - 2y^\alpha \frac{\partial}{\partial x} + 2x^\alpha \frac{\partial}{\partial y} - 2t^\alpha \frac{\partial}{\partial z}.$$

iii) Left-invariant Reeb fields  $\xi_1, \xi_2, \xi_3$  are  $\xi_1 = 2 \frac{\partial}{\partial x}$ ,  $\xi_2 = 2 \frac{\partial}{\partial y}$ ,  $\xi_3 = 2 \frac{\partial}{\partial z}$ .

iv) On  $\mathbf{G}(\mathbb{H})$ , the left-invariant connection is the Biquard connection. It is flat!

- Let  $\Psi \in \text{End}(H)$ . The  $Sp(n)$ -invariant parts are follows

$$\Psi = \Psi^{+++} + \Psi^{+--} + \Psi^{-+-} + \Psi^{--+}.$$

- The two  $Sp(n)Sp(1)$ -invariant components are given by

$$\Psi_{[3]} = \Psi^{+++}, \quad \Psi_{[-1]} = \Psi^{+--} + \Psi^{-+-} + \Psi^{--+}.$$

Using  $\text{End}(H) \stackrel{g}{\cong} \Lambda^{1,1}$  the  $Sp(n)Sp(1)$ -invariant components are the projections on the eigenspaces of  $\Upsilon = I_1 \otimes I_1 + I_2 \otimes I_2 + I_3 \otimes I_3$ .

## Quaternionic Contact Structure $(M^{4n+3}, \eta)$

- i) co-dim three distribution  $H$ , locally,  $H = \bigcap_{s=1}^3 \text{Ker } \eta_s$ ,  $\eta_s \in T^*M$ .
- ii)  $H$  carries a quaternion structure: a 2-sphere bundle of "almost complex structures" (locally) generated by  $I_s : H \rightarrow H$ ,  $I_s^2 = -1$ , satisfying  $I_1 I_2 = -I_2 I_1 = I_3$ ;
- iii) a "horizontal metric"  $g$  on  $H$ , such that for all  $X, Y \in H$

$$g(I_s X, I_s Y) = g(X, Y) \quad 2\omega_s(X, Y) \stackrel{\text{def}}{=} 2g(I_s X, Y) = d\eta_s(X, Y).$$

**Reeb vector fields:**  $TM = H \oplus V$ , for  $V = \text{span}\{\xi_1, \xi_2, \xi_3\}$  where

$$\eta_s(\xi_k) = \delta_{sk}, \quad (\xi_s \lrcorner d\eta_s)|_H = 0, \quad (\xi_s \lrcorner d\eta_k)|_H = -(\xi_k \lrcorner d\eta_s)|_H.$$

If  $n = 1$ , assume that the Reeb vector fields exist [**Duchemin, D.**].

**The Biquard connection:** There exists a unique linear connection  $\nabla$  on  $M$  with the properties: (i)  $V$  and  $H$  are parallel; (ii)  $g$  and

$\Omega = \sum_{j=1}^3 \omega_j \wedge \omega_j$  are parallel; (iii) the torsion satisfies

$$\blacktriangleright \forall X, Y \in H, \quad T(X, Y) = -[X, Y]|_V = 2\omega_i(X, Y)\xi_i \in V$$

$$\blacktriangleright \forall \xi \in V, X \in H, T_\xi(X) \equiv T(\xi, X) \in H \text{ and}$$

$$T_\xi \in (\mathfrak{sp}(n) + \mathfrak{sp}(1))^\perp, T_{\xi_j} = T_{\xi_j}^0 + I_j U, U \in \Psi_{[3]}.$$

$T_{\xi_j}^0$ -symmetric,  $I_j U$ -skew-symmetric..



We extend the horizontal metric  $g$  to a Riemannian metric  $h$  on  $M$  by requiring  $\text{span}\{\xi_1, \xi_2, \xi_3\} = V \perp H$  and  $h(\xi_s, \xi_t) = \delta_{st}$ .

**N.B.**  $h$  as well as the Biquard connection do not depend on the action of  $SO(3)$  on  $V$ .

- ▶ qc-curvature:  $\mathcal{R}(A, B)C = [\nabla_A, \nabla_B]C - \nabla_{[A, B]}C$
- ▶ qc-Ricci tensor:  $\text{Ric}(A, B) = \mathcal{R}(e_a, A, B, e_a) \stackrel{\text{def}}{=} \sum_{a=1}^{4n} h(\mathcal{R}(e_a, A)B, e_a)$ ;
- ▶ qc-scalar curvature:  $\text{Scal} = \text{tr}_H \text{Ric} = \text{Ric}(e_a, e_a)$ ;

**Theorem 13 (w/ Ivanov & Minchev '14).** If  $T^0 \stackrel{\text{def}}{=} T_{\xi_i}^0 I_i$ , then  $T^0 \in \Psi_{[-1]}$  and  $\text{Ric} = (2n + 2)T^0 + (4n + 10)U + \frac{\text{Scal}}{4n}g$ .

- ▶  $M$  is called **qc-Einstein** if  $T^0 = U = 0$ . For a qc-Einstein  $\Rightarrow \text{Scal} = \text{const}$  [w/ Ivanov & Minchev '10 & '1?] (non-trivial in 7-D, use  $W^{qc}$ !).  $M$  is called **qc-pseudo-Einstein** if  $U = 0$ .

**Theorem 14 (w/ Ivanov & Minchev, '14).** Suppose  $\text{Scal} > 0$ . The next conditions are equivalent:

- i)  $(M^{4n+3}, \eta)$  is qc-Einstein manifold.
- ii)  $M$  is locally 3-Sasakian

# Embedded qc manifolds [w/ Ivanov & Minchev arXiv:1406.4256]

**Theorem 15.** *If  $M$  is a qc-manifold embedded as a hypersurface in a hyper-Kähler manifold, then  $M$  is qc-conformal to a qc-Einstein structure. In particular, the qc Yamabe problem has a solution.*

**Theorem 16.** *If  $M$  is a connected qc-hypersurface of  $\mathbb{R}^{4n+4} \cong \mathbb{H}^{n+1}$  then, up to a quaternionic affine transformation of  $\mathbb{H}^{n+1}$ ,  $M$  is contained in one of the following three hyperquadrics:*

$$(i) |q_1|^2 + \cdots + |q_n|^2 + |p|^2 = 1, \quad (ii) |q_1|^2 + \cdots + |q_n|^2 - |p|^2 = -1, \\ (iii) |q_1|^2 + \cdots + |q_n|^2 + \operatorname{Re}(p) = 0.$$

*Here  $(q_1, q_2, \dots, q_n, p)$  denote the standard quaternionic coordinates of  $\mathbb{H}^{n+1}$ . In particular, if  $M$  is a compact qc-hypersurface of  $\mathbb{R}^{4n+4} \cong \mathbb{H}^{n+1}$  then, up to a quaternionic affine transformation of  $\mathbb{H}^{n+1}$ ,  $M$  is the standard 3-Sasakian sphere.*

## Standard qc-structure on 3-Sasakian sphere

- ▶ Contact 3-form on the sphere  $S = \{|q|^2 + |p|^2 = 1\} \subset \mathbb{H}^n \times \mathbb{H}$ ,

$$\tilde{\eta} = dq \cdot \bar{q} + dp \cdot \bar{p} - q \cdot d\bar{q} - p \cdot d\bar{p}.$$

- ▶ Identify  $\mathbf{G}(\mathbb{H})$  with the boundary  $\Sigma$  of a Siegel domain in  $\mathbb{H}^n \times \mathbb{H}$ ,

$$\Sigma = \{(q', p') \in \mathbb{H}^n \times \mathbb{H} : \operatorname{Re} p' = |q'|^2\},$$

by using the map  $(q', \omega') \mapsto (q', |q'|^2 - \omega')$ .

**Proposition 17.** *The Cayley transform,  $\mathcal{C} : S \setminus \{pt.\} \rightarrow \Sigma$ ,*

$$(q', p') = \mathcal{C} \left( (q, p) \right) = ((1 + p)^{-1} q, (1 + p)^{-1} (1 - p)).$$

*is a qc-conformal transformation*

$$\mathcal{C}^* \tilde{\Theta} = \frac{1}{2|1 + p|^2} \lambda \tilde{\eta} \bar{\lambda}, \quad \lambda - \text{unit quaternion.}$$

## QC divergence formula

**Theorem 18 (w/ Ivanov & Minchev arXiv:1504.03142).** Suppose  $(M^{4n+3}, \eta)$  is a qc structure which is qc-conformal to a qc-Einstein structure  $(M^{4n+3}, \bar{\eta})$ ,  $\bar{\eta} = \frac{1}{2h} \eta$ . If  $Scal_{\eta} = Scal_{\bar{\eta}} = 16n(n+2)$ , then  $(M^{4n+3}, \eta)$  is also qc-Einstein. In fact, with  $f = \frac{1}{2} + h + \frac{1}{4}h^{-1}|\nabla h|^2$ , we have

$$\begin{aligned} \nabla^* \left( f(D + E) + \sum_{s=1}^3 dh(\xi_s) \left( I_s E + F_s + 4I_s A_s - \frac{10}{3} I_s A \right) \right) \\ = \left( \frac{1}{2} + h \right) \left( |T^0|^2 + 4|U|^2 \right) + 2h|\mathbb{D} + \mathbb{E}|^2 + h \langle QV, V \rangle. \end{aligned}$$

where  $Q$  is a positive definite matrix,  $V = (E, D_1, D_2, D_3, A_1, A_2, A_3)$ , and

$$\begin{aligned} E = -2h^{-1}U\nabla h, \quad D_i = -\frac{1}{2}h^{-1}(T^0 - I_i T^0 I_i)\nabla h, \quad F_i = -\frac{1}{2}h^{-1}T^0 I_i \nabla h, \\ A_i = I_i[\xi_j, \xi_k], \quad A = \sum_{i=1}^3 A_i, \quad D = -h^{-1}T^0 \nabla h. \end{aligned}$$

## Infinitesimal QC transformations w/ Ivanov & Minchev '14

A vector field  $Q$  on a qc manifold  $(M, \eta)$  is a *qc vector field* if its flow preserves the horizontal distribution  $H = \ker \eta$ ,

$$\mathcal{L}_Q \eta = (\nu I + O) \cdot \eta,$$

where  $\nu$  is a smooth function and  $O \in so(3)$  with smooth entries.

Thus, we also have

$$\mathcal{L}_Q g = \nu g, \quad \mathcal{L}_Q I = O \cdot I, \quad I = (I_1, I_2, I_3)^t.$$

The function  $\nu = \frac{1}{2n} \nabla^* Q_H$  since

$$g(\nabla_X Q_H, Y) + g(\nabla_Y Q_H, X) + 2\eta_s(Q)g(T_{\xi_s}^0 X, Y) = \nu g(X, Y).$$

The infinitesimal version of the qc Yamabe equation for a qc vector field is

**Proposition 19.** *Let  $(M^{4n+3}, \eta)$  be a qc manifold. For any qc vector field  $Q$  on  $M$  we have*

$$\Delta(\nabla^* Q_H) = -\frac{n}{2(n+2)} Q(\text{Scal}) - \frac{\text{Scal}}{4(n+2)} \nabla^* Q_H.$$

**Lemma 20.** *Let  $(M, \eta)$  and  $(M, \bar{\eta})$  be qc-Einstein manifolds with equal qc-scalar curvatures  $16n(n+2)$ . If  $\bar{\eta} = \frac{1}{2h}\eta$  for some smooth  $h > 0$ , then*

$$Q = \frac{1}{2}\nabla f + \sum_{s=1}^3 dh(\xi_s)\xi_s$$

*is a qc vector field on  $M$ , where  $f = \frac{1}{2} + h + \frac{1}{4}h^{-1}|\nabla h|^2$  is the function in the divergence formula.*

It follows,  $\phi = \frac{1}{2}\Delta f$  is an eigenfunction of the sub-Laplacian with eigenvalue  $-4n$  unless  $\Delta f \equiv 0$ . In the first case, the qc version of the Lichnerowicz-Obata eigenfunction sphere theorem shows that  $(M, \eta)$  is the 3-Sasakain sphere. If  $\Delta f = 0$ , then  $f = \frac{1}{2} + h + \frac{1}{4}h^{-1}|\nabla h|^2 = \text{const}$  since  $M$  is compact. It follows that  $h = 1/2$  by considering the points where  $h$  achieves its minimum and maximum and using the qc Yamabe equation.

## QC Lichnerowicz

**Theorem 21 (w/ Ivanov, S., & Petkov, A. '13 & '14).** *Let  $(M, \eta)$  be a compact QC manifold of dimension  $4n + 3$ . Suppose, for  $\alpha_n = \frac{2(2n+3)}{2n+1}$ ,  $\beta_n = \frac{4(2n-1)(n+2)}{(2n+1)(n-1)}$  and for any  $X \in H$*

$$\mathcal{L}(X, X) \stackrel{\text{def}}{=} 2Sg(X, X) + \alpha_n T^0(X, X) + \beta_n U(X, X) \geq 4g(X, X).$$

*If  $n = 1$ , assume in addition the positivity of the P-function of any eigenfunction. Then, any eigenvalue  $\lambda$  of the sub-Laplacian  $\Delta$  satisfies the inequality*

$$\lambda \geq 4n$$

The 3-Sasakian sphere achieves equality in the Theorem. The eigenspace of the first non-zero eigenvalue of the sub-Laplacian on the unit 3-Sasakian sphere in Euclidean space is given by the restrictions to the sphere of all linear functions.

## Definition of the QC P-function

a) The  $P$ -form of a function  $f$  is the 1-form

$$P_f(X) = \nabla^3 f(X, e_b, e_b) + \sum_{t=1}^3 \nabla^3 f(I_t X, e_b, I_t e_b) \\ - 4n Sdf(X) + 4n T^0(X, \nabla f) - \frac{8n(n-2)}{n-1} U(X, \nabla f).$$

b) The  $P$ -function of  $f$  is the function  $P_f(\nabla f)$ .

c) The  $C$ -operator is the 4-th order differential operator on  $M$  (independent of  $f$ !)

$$f \mapsto Cf = \nabla^* P_f = (\nabla_{e_a} P_f)(e_a).$$

d) The  $P$ -function of  $f$  is non-negative if

$$\int_M f \cdot Cf \operatorname{Vol}_\eta = - \int_M P_f(\nabla f) \operatorname{Vol}_\eta \geq 0.$$

If the above holds for any  $f \in \mathcal{C}_o^\infty(M)$  we say that the  $C$ -operator is non-negative,  $C \geq 0$ .



## Properties of the C-operator

**Theorem 22 (w/ Ivanov & Petkov, '13).** a)  $C \geq 0$  for  $n > 1$ .

Furthermore  $Cf = 0$  iff  $(\nabla^2 f)_{[3][0]}(X, Y) = 0$ . In this case the  $P$ -form of  $f$  vanishes as well.

b) If  $n = 1$  and  $M$  is qc-Einstein with  $Scal \geq 0$ , the  $P$ -function of an eigenfunction of the sub-Laplacian is non-negative,

$$\Delta f = \lambda f \quad \Rightarrow \quad - \int_M P_f(\nabla f) \text{Vol}_\eta \geq 0.$$

- ▶  $(\nabla_{e_a}(\nabla^2 f)_{[3][0]})(e_a, X) = \frac{n-1}{4n} P_f(X)$ , hence

$$\frac{n-1}{4n} \int_M f \cdot Cf \text{Vol}_\eta = -\frac{n-1}{4n} \int_M P_f(\nabla f) \text{Vol}_\eta = \int_M |(\nabla^2 f)_{[3][0]}|^2 \text{Vol}_\eta,$$

after using the Ricci identities, the divergence formula and the orthogonality of the components of the horizontal Hessian.

- ▶ qc-Einstein  $\Rightarrow Scal = const$ ,  $\nabla^3 f(\xi_s, X, Y) = \nabla^3 f(X, Y, \xi_s)$ , and the vertical space is integrable;  $\nabla^2 f(\xi_k, \xi_j) - \nabla^2 f(\xi_j, \xi_k) = -Sdf(\xi_i)$
- ▶  $\int_M |P_f|^2 \text{Vol}_\theta = -(\lambda + 4S) \int_M P_f(\nabla f) \text{Vol}_\theta$

## The QC Obata type theorem in the compact case

**Theorem 23 (w/ Ivanov & Petkov, arxiv1303.0409).** *Let  $(M, \eta)$  be a compact QC manifold of dimension  $4n + 3$  which satisfies a Lichnerowicz' type bound  $\mathcal{L}(X, X) \geq 4g(X, X)$ . Then, there is a function  $f$  with  $\Delta f = 4nf$  if and only if  $M$  is qc-homothetic to the 3-Sasakian sphere, assuming in addition  $M$  is qc-Einstein when  $n = 1$ .*

Remarks:

- ▶ The 7-D case is still open in the general case.
- ▶ The results follow from another theorem where only completeness and knowledge of the horizontal Hessian are assumed.

# Proof of QC eigenvalue Obata for a qc-Einstein

1. Show that  $(\nabla^h)^2 f(X, Y) = -fh(X, Y)$ , ( $h$ - Riemannian metric!).
2. Obata's result shows  $(M, h)$  is homothetic to the unit sphere in quaternion space.
3. Show qc-conformal flatness.
4. Use the qc-Liouville theorem to see  $(M, g, \eta, \mathbb{Q})$  is qc-conformal to  $S^{4n+3}$ , i.e., we have  $\eta = \kappa \Psi F^* \tilde{\eta}$  for some diffeomorphism  $F : M \rightarrow S^{4n+3}$ ,  $0 < \kappa \in \mathcal{C}^\infty(M)$ , and  $\Psi \in \mathcal{C}^\infty(M : SO(3))$

**Theorem 24 (Čap, A., & Slovák, J., '09; w/ Ivanov & Petkov arXiv:1303.0409).** Every qc-conformal transformation between open subsets of the 3-Sasakian unit sphere is the restriction of a global qc-conformal transformation.

Rmrk: Cowling, M., & Ottazzi, A., Conformal maps of Carnot groups, arXiv:1312.6423.

**Theorem 25 (w/ Ivanov & Minchev).** Let  $\Theta = \frac{1}{2h} \tilde{\Theta}$  be a conformal deformation of the standard qc-structure  $\tilde{\Theta}$  on the quaternionic Heisenberg group  $\mathbf{G} (\mathbb{H})$ . If  $\Theta$  is also qc-Einstein, then

$$h(q, \omega) = c_0 \left[ (\sigma + |q + q_0|^2)^2 + |\omega + \omega_0 + 2 \operatorname{Im} q_0 \bar{q}|^2 \right].$$

with  $c_0 > 0$  and  $\sigma \in \mathbb{R}$ . Furthermore,  $S_\Theta = 128n(n+2)c_0\sigma$ .

5. compare the metrics on  $H$  to see homothety.

## QC Conformal Curvature tensor

- ▶ "Schouten" tensor  $L(X, Y) = \frac{1}{2}T^0(X, Y) + U(X, Y) + \frac{Scal}{32n(n+2)} g(X, Y)$ .
- ▶ Conformal curvature

$$\begin{aligned}W^{qc}(X, Y, Z, V) &= R(X, Y, Z, V) + (g \otimes L)(X, Y, Z, V) \\&\quad + \sum_{s=1}^3 (\omega_s \otimes I_s L)(X, Y, Z, V) \\&\quad - \frac{1}{2} \sum_{(i,j,k)} \omega_i(X, Y) \left[ L(Z, I_i V) - L(I_i Z, V) + L(I_j Z, I_k V) - L(I_k Z, I_j V) \right] \\&\quad - \sum_{s=1}^3 \omega_s(Z, V) \left[ L(X, I_s Y) - L(I_s X, Y) \right] + \frac{1}{2n} (trL) \sum_{s=1}^3 \omega_s(X, Y) \omega_s(Z, V),\end{aligned}$$

where  $\sum_{(i,j,k)}$  denotes the cyclic sum.

$W^{qc}$  is qc-conformal invariant, i.e., if  $\bar{\eta} = \kappa \Psi \eta$  then  $W_{\bar{\eta}}^{qc} = \phi W_{\eta}^{qc}$ ,  $0 < \kappa \in \mathcal{C}^{\infty}(M)$ , and  $\Psi \in \mathcal{C}^{\infty}(M : SO(3))$

**Theorem 26 (w/ Ivanov '10).** A qc manifold is locally qc-conformal to the quaternionic sphere  $S^{4n+3}$  or quaternion Heisenberg group iff the qc conformal curvature vanishes,  $W^{qc} = 0$ .