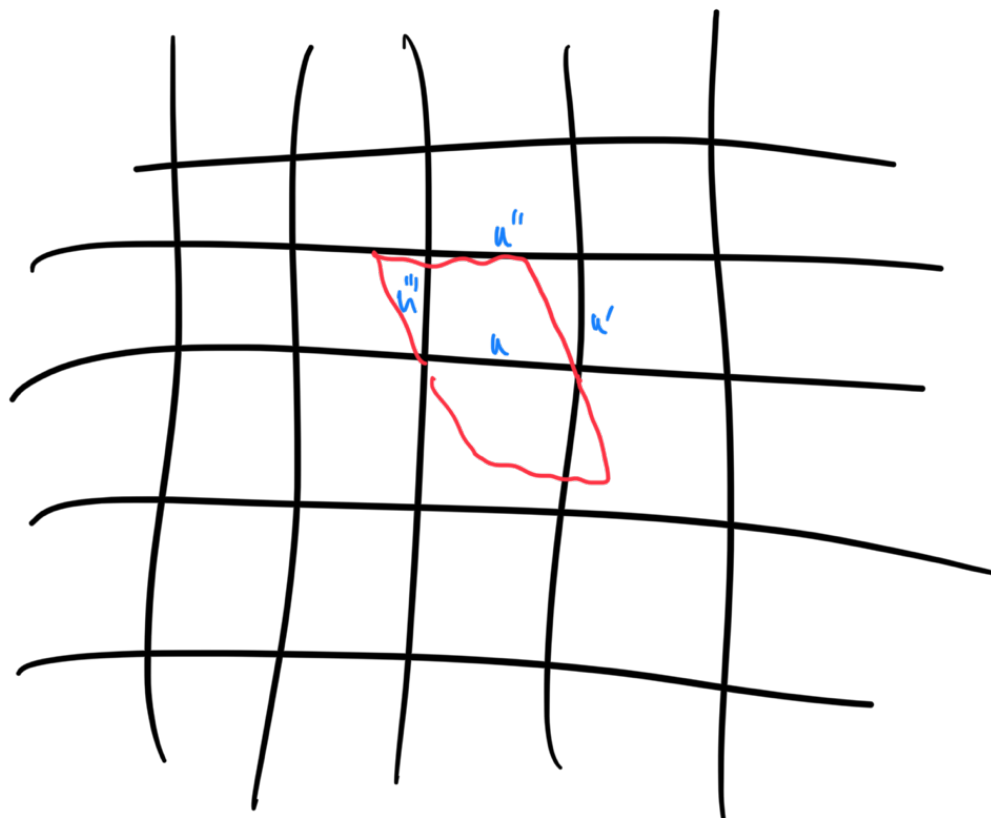




$d=4$ 10^4 lattice points

$100^4 = 10^8$ points

$$\langle A \rangle = \frac{\int e^{-S_E(\phi)} A(\phi) D\phi}{\int e^{-S_E(\phi)} D\phi}$$

$$ds^2 = g_{ik} dx^i dx^k$$

$$g_{ik} \approx \eta_{ik} =$$

$$\begin{pmatrix} -1 & & \\ & 1 & \\ & & 1 \end{pmatrix}$$

$$A^i{}_e B^e{}_j \equiv \sum_{e=0}^3 A^i{}_e B^e{}_j$$

$$L \eta L = \eta$$

$$\eta^2 = 1$$

$$dx'^k = L^k{}_e dx^e$$

$$A'^K = \mathcal{L}^K_{\ell} A$$

contra

$$\frac{\partial}{\partial x'^i} = \frac{\partial x^\ell}{\partial x'^i} \frac{\partial}{\partial x^\ell}$$



$$\int_{\partial R} \mathbf{E} \cdot d\mathbf{r} = \int_R \nabla \times \mathbf{E} \cdot d\mathbf{S}$$

boundary of



$$(r^2 R'_{k\ell}(r))' + [h^2 r^2 - \ell(\ell+1)] R_{k\ell}(r) = 0$$

$$R_{k\ell}(r) = j_\ell(hr)$$

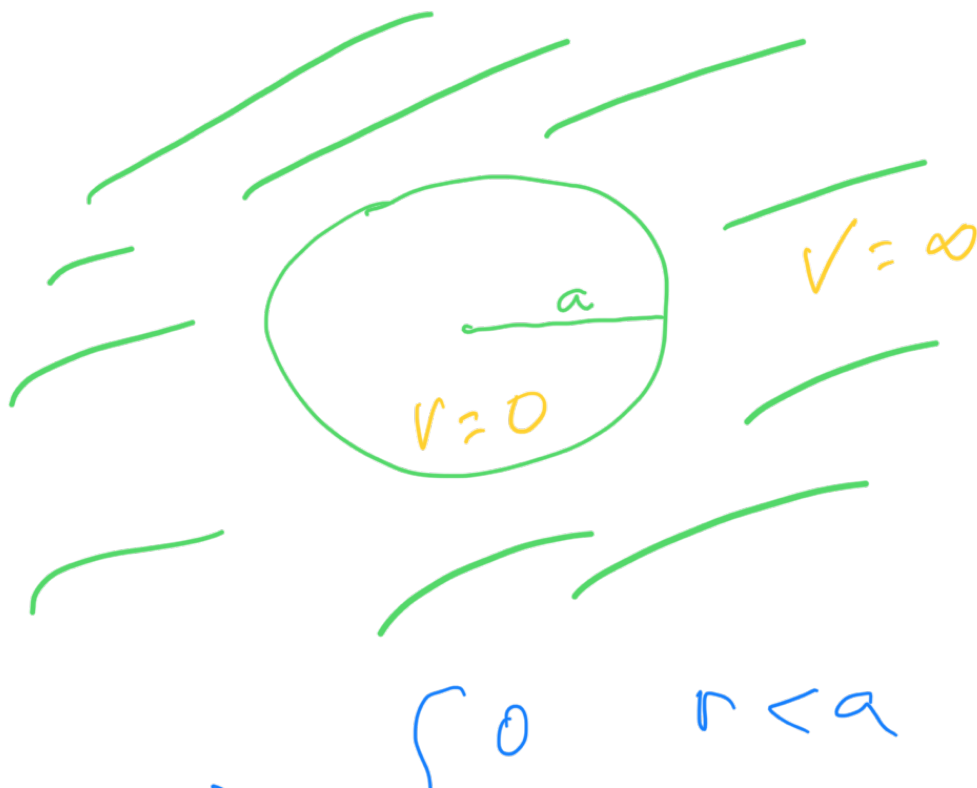
$$\frac{d}{dr} \left(r^2 \frac{d}{dr} j_\ell(hr) \right) + [h^2 r^2 - \ell(\ell+1)] j_\ell(hr) = 0$$

$$k \frac{d}{dkh} \left(r^2 k \frac{d}{dkh} j_\ell(hr) \right) + [h^2 r^2 - \ell(\ell+1)] j_\ell(hr) = 0$$

$x = hr$

$$\frac{d}{dx} \left(x^2 \frac{d}{dx} j_\ell(x) \right) + [x^2 - \ell(\ell+1)] j_\ell(x) = 0$$

$-\Delta V = h^2 V$



$$V(r) = \begin{cases} \infty & r > a \end{cases}$$

$$\text{So } j_l(ka) = 0 = R_{kl}(a)$$

$$\text{So } ka = z_{l,n} \quad j_l(z_{l,n}) = 0$$

$$-\frac{\hbar^2}{2m} \nabla^2 \psi_{lm} = E \psi_{lm}$$

$$+\nabla^2 \psi_{lm} = -\frac{2mE}{\hbar^2} \psi_{lm} = -k^2 \psi_{lm}$$

$$k^2 = \frac{2mE}{\hbar^2} = \left(\frac{z_{l,n}}{a}\right)^2$$

$$E_{lm} = \frac{\hbar^2}{2m} \left(\frac{z_{l,n}}{a}\right)^2$$

So, the smaller the size of the volume in which the particle is free to move, the higher its energy.

$$E = h\nu$$

$$x = (x^0, x^1, x^2, x^3) \quad \text{eg. } x = (ct, \vec{r})$$

$$dx^k \quad dx^1 = x^1 + \epsilon - x^1 = \epsilon$$

$$dx'^i = \frac{\partial x'^i}{\partial x^k} dx^k$$

$$\begin{pmatrix} x'^0 \\ x'^1 \\ x'^2 \\ x'^3 \end{pmatrix} = \begin{pmatrix} \frac{\partial x'^0}{\partial x^0} & \dots & \frac{\partial x'^0}{\partial x^3} \\ \vdots & \ddots & \vdots \\ \frac{\partial x'^3}{\partial x^0} & \dots & \frac{\partial x'^3}{\partial x^3} \end{pmatrix} \begin{pmatrix} x^0 \\ x^1 \\ x^2 \\ x^3 \end{pmatrix}$$

contra variant
vectors

$$A'^k(x') = \sum_{i=0}^3 \frac{\partial x'^k}{\partial x^i} A^i(x)$$

$x \rightarrow x'$

field = something that's
a function of $x = (t, \vec{x})$

$$\frac{\partial}{\partial x''^i} = \sum_{k=0}^3 \frac{\partial x^k}{\partial x''^i} \frac{\partial}{\partial x^k}$$

chain rule

covariant vectors

$$\rightarrow C'_i(x') = \sum_{k=0}^3 \frac{\partial x^k}{\partial x''^i} C_k(x)$$

covariant vector field,

$$F'_{je}(x') = \sum_{ik=0}^3 \frac{\partial x^i}{\partial x''^j} \frac{\partial x^k}{\partial x''^e} F_{ik}(x)$$

correlation tensor
(of rank 2)

$$A_i B^i \equiv \sum_{i=0}^3 A_i B^i$$

$\langle 4 | S | 4 \rangle =$ average
value of S in state Ψ .

$$e_i(x) \cdot e_k(x) = f_{ik}(x) \\ = \sum_{\alpha=1}^D e_i^{\alpha}(x) e_k^{\alpha}(x)$$

empty space

$$n(x) = n_{ir}(x)$$

$$g_{ik} = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & a & 0 & 0 \\ 0 & 0 & a & 0 \\ 0 & 0 & 0 & a \end{pmatrix}$$

$$\frac{\partial V}{\partial x^k} = \frac{\partial (v^i e_i)}{\partial x^k} = \frac{\partial v^i}{\partial x^k} e_i + v^i \frac{\partial e_i}{\partial x^k}$$

covariant derivative

$$D_e V^k = e^k \cdot \frac{\partial V}{\partial x^k}$$

$$e_i \cdot e_k = g_{ik}$$

$$e^i \cdot e_k = \delta^i_k$$

$$D_e V^k = e^k \cdot \left(\frac{\partial v^i}{\partial x^k} e_i + v^i \frac{\partial e_i}{\partial x^k} \right)$$

$$= \delta^k_i \frac{\partial v^i}{\partial x^k} + e^k \cdot \frac{\partial e_i}{\partial x^k} v^i$$

$$= v^k \quad - k \frac{\partial e_i}{\partial x^k} v^i$$

$$= \frac{\partial v}{\partial x^l} + \underbrace{e_i \cdot \frac{\partial e^i}{\partial x^l}}_{\Gamma_{il}^k} V^i$$

$$D_l V^k = \frac{\partial V^k}{\partial x^l} + \underbrace{\Gamma_{il}^k}_{\text{red line}} V^i$$

$$V = V_k e^k$$

$$\frac{\partial v}{\partial x^l} = \frac{\partial V_k}{\partial x^l} e^k + V_k \frac{\partial e^k}{\partial x^l}$$

$$e_i \cdot \frac{\partial v}{\partial x^l} = \frac{\partial V_k}{\partial x^l} e_i \cdot e^k + V_k e_i \cdot \frac{\partial e^k}{\partial x^l}$$

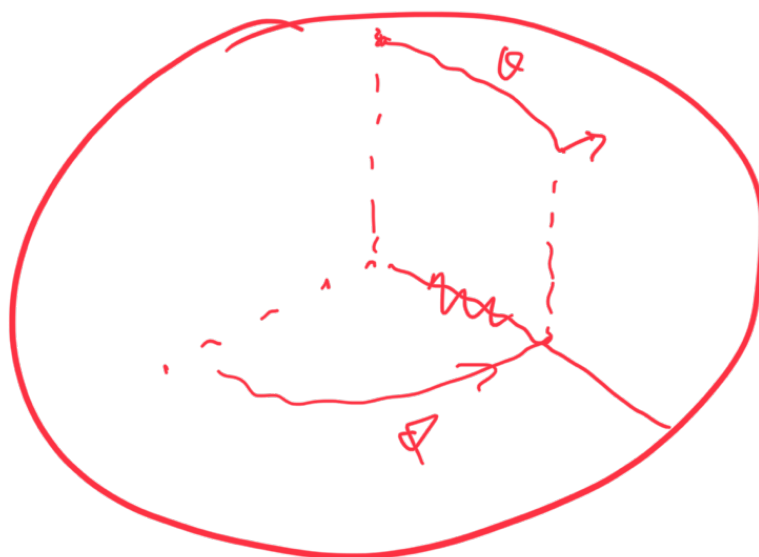
$$= \frac{\partial V_i}{\partial x^l} + V_k \underbrace{e_i \cdot \frac{\partial e^k}{\partial x^l}}_{\Gamma_{il}^k}$$

$$D_l V_i = \frac{\partial V_i}{\partial x^l} + V_k \underbrace{e_i \cdot \frac{\partial e^k}{\partial x^l}}_{\Gamma_{il}^k}$$

$$= \frac{\partial V_i}{\partial x^k} - \Gamma^i_{jk} V^j$$

covariant derivative
of a covariant vector.

$$D_x \eta_{ik} = 0$$



$$d\phi = \hat{\phi} r \sin\theta d\phi$$

in 3D

$$e_i(x) = \partial_i \phi(x)$$

$$e_\theta(x) = \partial_\theta p(x)$$

$$e_\phi(x) = \partial_\phi p(x)$$

complete on
surface

but they are vectors in
3D

e_i basis vectors.

$$g_{ik}(x) = e_i(x) \cdot e_k(x)$$

$$g^{li} g_{ik} = \delta^l_k$$

$$A^i(x) = g^{ik} A_k(x)$$

$$A_k(x) = g_{ke} A^e(x)$$

$$e^i = g^{ik} e_k$$

...

$$\begin{aligned}
 V(x) &= V^i(x) e_i(x) \\
 V'(x') &= V'^i(x') e'_i(x') \\
 &= \frac{\partial x'^i}{\partial x^j} V^j(x) \frac{\partial x^k}{\partial x'^i} e_k(x) \\
 &= \int_j^k V^j(x) e_k(x) \\
 &= V^i(x) e_i(x) = V(x)
 \end{aligned}$$

So $V(x)$ is a scalar.

$$\frac{\partial V}{\partial x^e} = \frac{\partial V^i}{\partial x^e} e_i(x) + V^i(x) \frac{\partial e_i(x)}{\partial x^e}$$

covariant vector

$$D_e V^i = \frac{\partial V^i}{\partial x^e} e_i(x) + \frac{\partial e_i(x)}{\partial x^e} V^i(x)$$

$$e^k \cdot D_e V^i = \frac{\partial V^k}{\partial x^e} + e^k \cdot \frac{\partial e_i(x)}{\partial x^e} V^i(x)$$

$$e^k \cdot e_i = \delta^k_i$$

$k, i = 1, \dots, n$

$$\Gamma^k_{il} = e^i \cdot \frac{\partial e^k}{\partial x^l}$$

$$D_x V^k = \partial_x V^k + \Gamma^k_{il} V^i$$

$$D_x \psi = \partial_x \psi + \underbrace{A_x \psi}_{\text{photon}}$$

$$D_x V_i = \partial_x V_i - \Gamma^k_{ie} V_k$$

$$e_i \cdot e_k = e_i^\alpha \eta_{\alpha\beta} e_k^\beta$$

$$\eta = \begin{pmatrix} -1 & & & \\ & -1 & & \\ & & +1 & \\ & & & +1 \end{pmatrix}$$

$$g_{ik} = e_i \cdot e_k = e_i^\alpha \eta_{\alpha\beta} e_k^\beta$$

$$A_{ie} = -A_{ei}$$

$$\partial_k A_{ie} + \partial_e A_{ki} + \partial_i A_{ek} \\ = \partial_k A_{ie} + \partial_e A_{ki} + \partial_i A_{ek}$$

$$F_{ik} = -F_{ki}$$

$$\Gamma^k_{il} = \frac{1}{2} g^{kj} (\partial_l g_{ji} + \partial_i g_{jl} - \partial_j g_{il})$$

$$ds^2 = g_{ik} dx^i dx^k$$

4D geometry →

$a(t)$: how space expands
 space is maximally symmetric

$$F_{..} = [\partial_{..} + A_i \partial_i + A_k \partial_k]$$

$$= \partial_i A_k - \partial_k A_i$$

$$+ A_i A_k - A_k A_i$$

$$g' + xy = c e^{-x^2/2}$$

$$r(x) = x$$

$$\int r(x) dx$$

$$d(x) = e$$

$$s(x) = c e^{-x^2/2}$$

$$= e^{x^2/2}$$

$$y(x) = g(0) e^{-x^2/2}$$

$$+ e^{-x^2/2} \int c dx$$

$$= (g(0) + cx) e^{-x^2/2}$$

$$y(0) = (g(0) + 0) e^{-0^2/2} = g(0).$$

$$V(r, \theta) = \sum \frac{b_l}{r^{l+1}} P_l(\cos \theta)$$

$$V(R, \theta) = \sum \frac{b_l}{R^{l+1}} P_l(\cos \theta)$$

$$= V_0 \cos^2 \theta$$

$$x^2 \Rightarrow \frac{2 P_2(x) + P_0(x)}{3}$$

$$\sum \frac{b_l}{R^{l+1}} P_l(x) = \frac{V_0}{3} (2 P_2(x) + P_0(x))$$

$$\int \sum \frac{b_l}{R^{l+1}} P_l(x) P_m(x) dx$$

$$= \sum \frac{b_l}{R^{l+1}} \frac{2}{2m+1} \delta_{ml}$$

$$= \frac{2}{2m+1} \frac{b_m}{R^{m+1}}$$

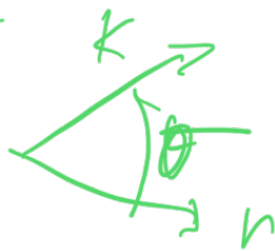
$$= \frac{V_0}{3} \left(2 \frac{2}{5} \delta_{m2} + 2 \delta_{m0} \right)$$

$$b_m = \frac{(2m+1) V_0 R^{m+1}}{3} \left(\frac{2}{5} \delta_{m2} + \delta_{m0} \right)$$

$$b_0 = \frac{V_0 R}{3} \quad b_2 = \frac{2 V_0 R^3}{3}$$

$$V(r, \theta) = \frac{V_0}{3} \left[\frac{R}{r} + \frac{R^3}{r^3} (3 \cos^2 \theta - 1) \right]$$

$$\underline{k \cdot n = kn \cos \theta}$$



$$k = |\vec{k}|$$

$$n = |\vec{n}|$$

$$i k \cdot n$$

$$\langle \vec{n} | \vec{k} \rangle = \frac{e^{i k \cdot n}}{(2\pi)^{3/2}}$$

$$= \frac{e^{i k n \cos \theta}}{(2\pi)^{3/2}} = \frac{1}{(2\pi)^{3/2}} \sum_{l=0}^{\infty} (2l+1) P_l(\cos \theta) i^l j_l(kn)$$

$$= \frac{1}{(2\pi)^{3/2}} \sum_{l=0}^{\infty} \sum_{m=-l}^l 4\pi i^l j_l(kn) Y_{lm}(\theta, \phi)$$

$$\theta \phi \leftrightarrow \vec{r}$$

$$r' \theta' \phi' \leftrightarrow \vec{r}'$$

